

# Lecture L05 - ASTC25

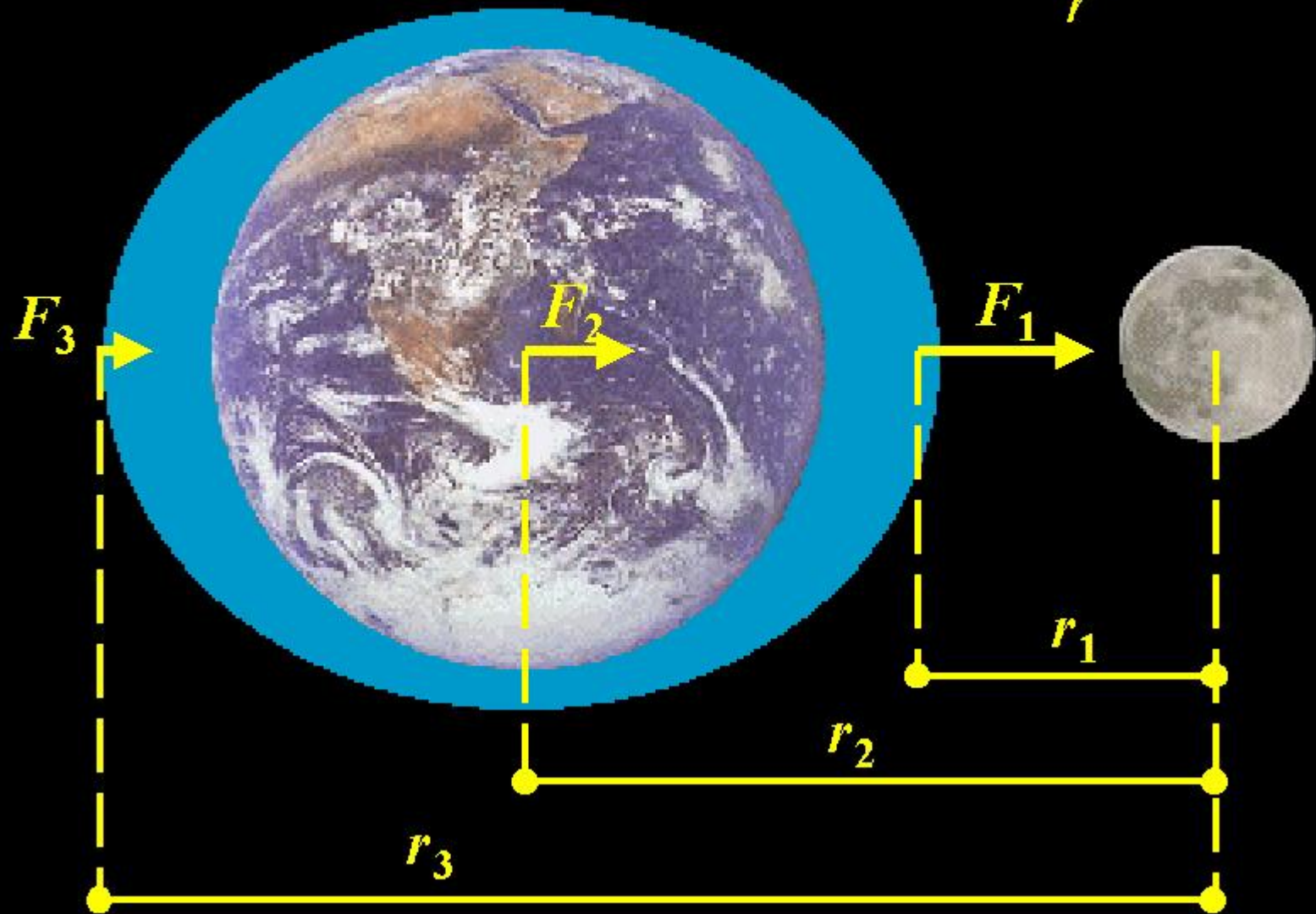
## Physics of Tides

Tides in the solar system

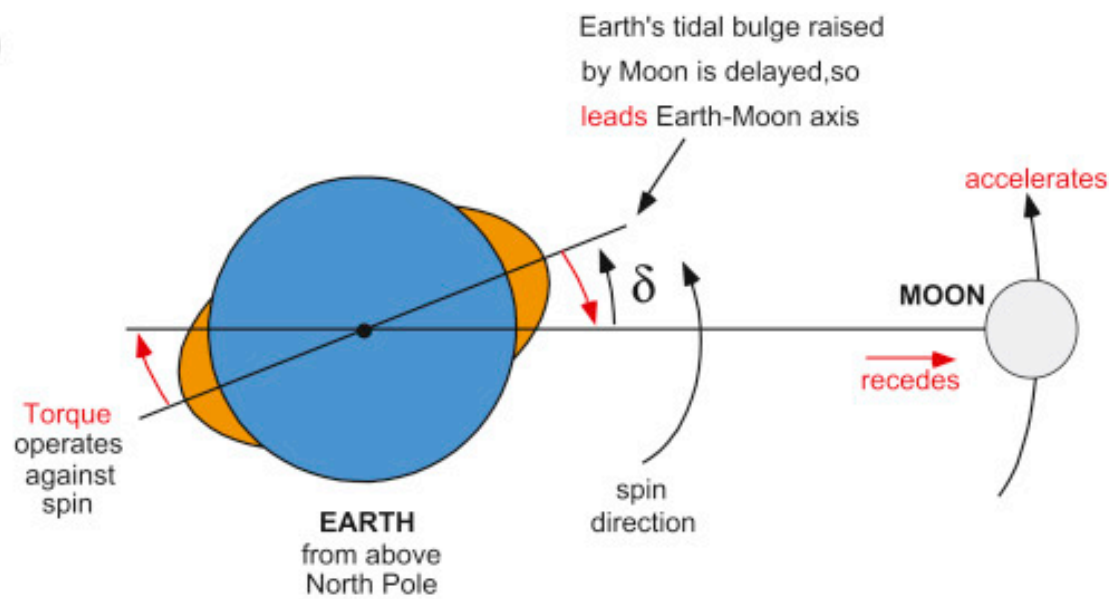
Roche limit of tidal disruption

Please also read Lissauer & de Pater textbook, ch. 2.7  
& see wikipedia article on “Volcanism on Io”

$$F = \frac{G M_{\text{moon}} m}{r^2}$$



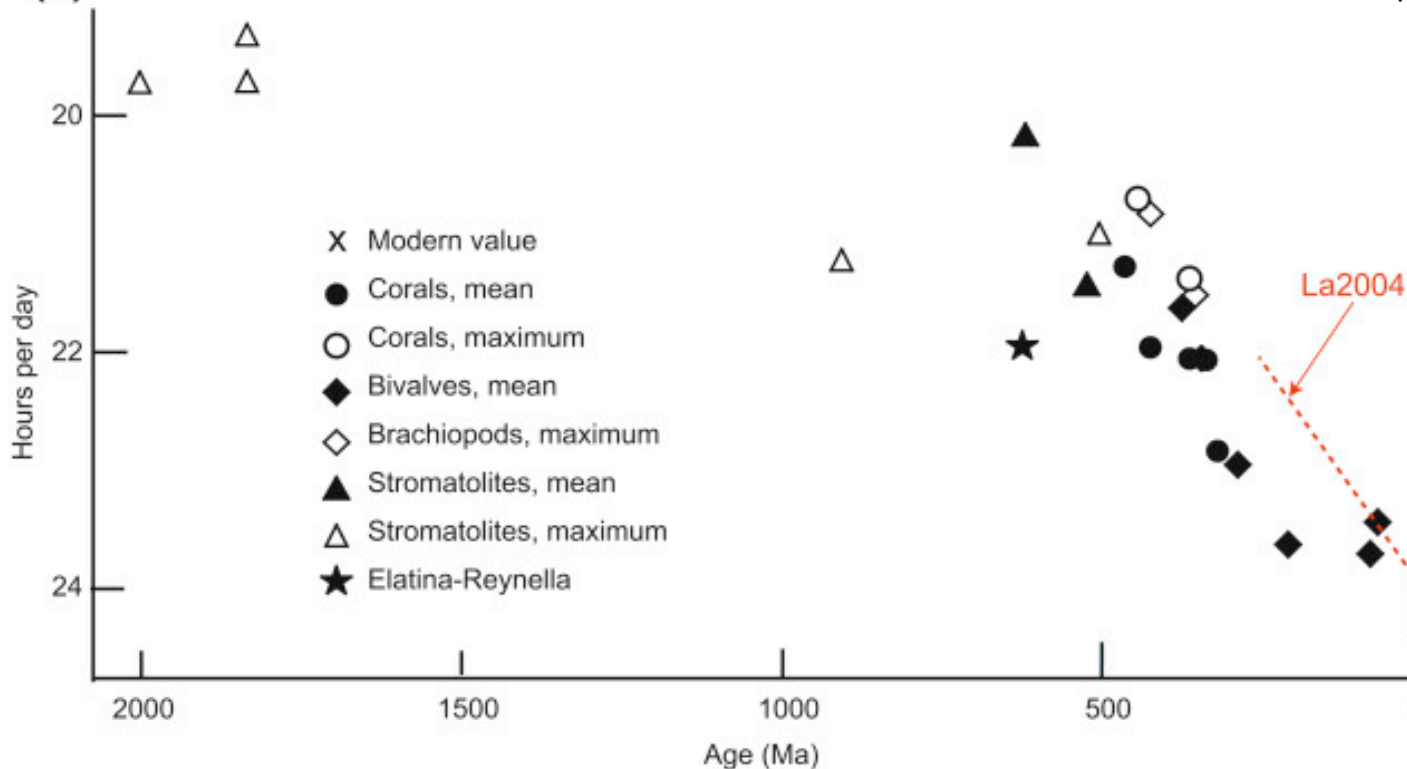
(a)



$\delta$  is called the tidal lag angle

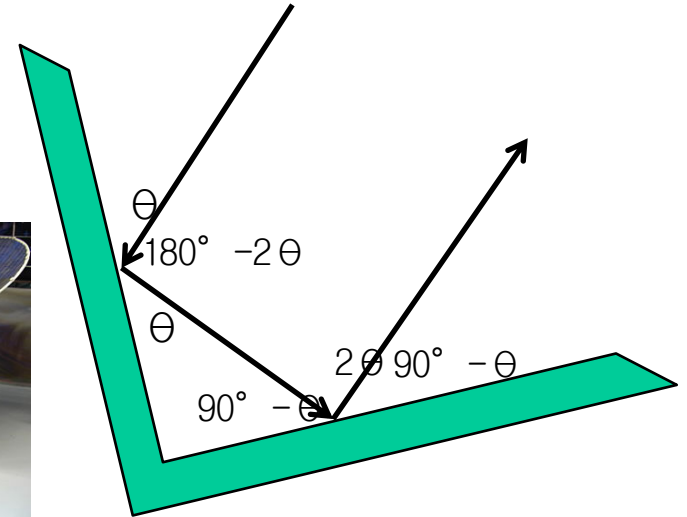
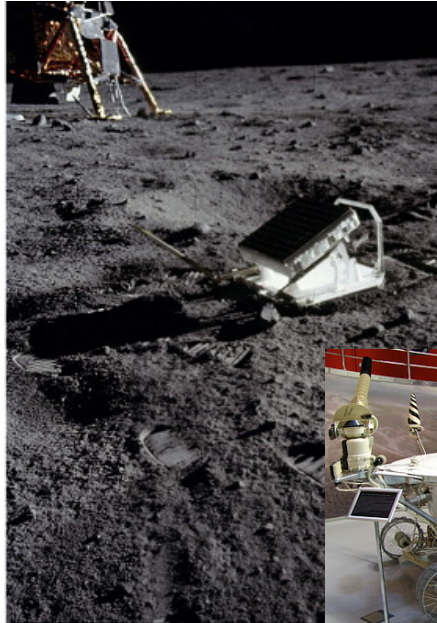
“accelerates” is an interesting misnomer: the Moon by being pulled forward actually travels slower & slower because of Kepler’s laws

(b)



As the Earth spins down, the day lengthens i.e. lasts more and more current constant hours

In 1970s **retro-reflectors** (mirrors) were left on the surface of the Moon by American Apollo (astronautic) and Soviet Lunokhod mission (1<sup>st</sup> robotic planetary rover). We use them to track the distance to the Moon with accuracy better than 2 cm.



The Moon's distance currently **increases 3.76 cm/yr**, and on longer

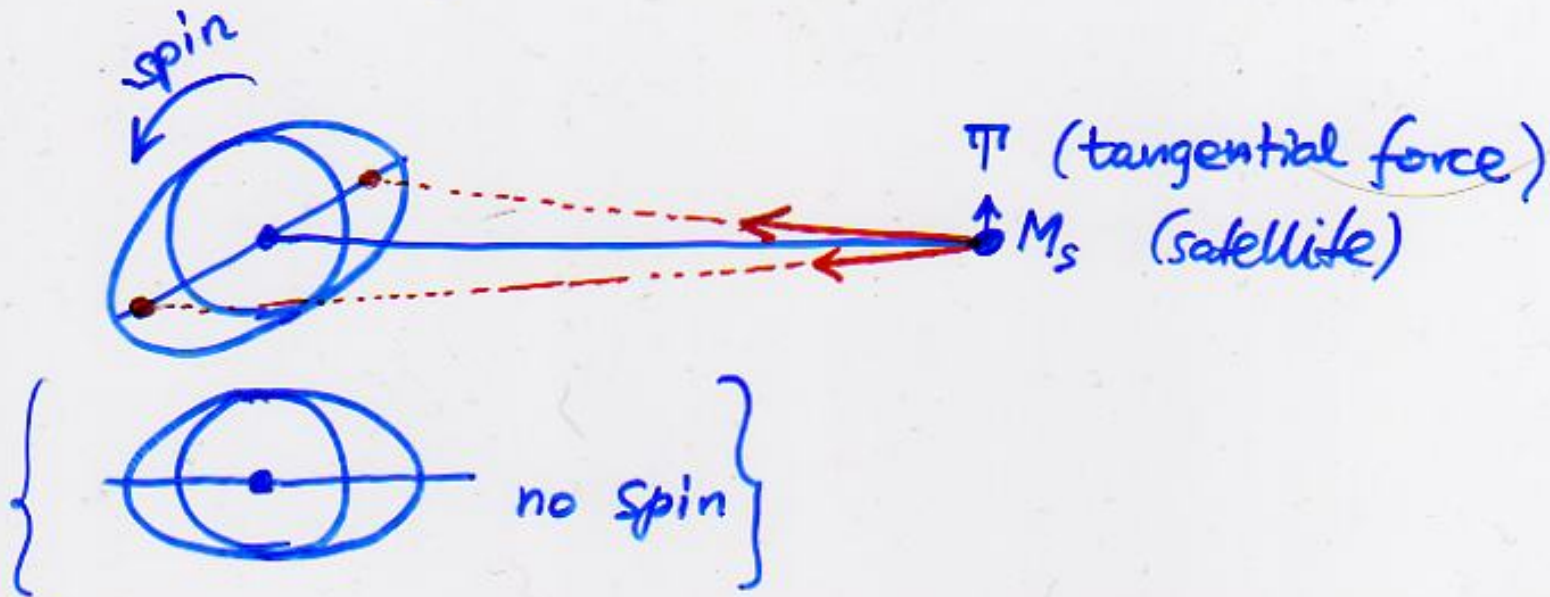
time scales the semi-major axis grows as follows:

$$a_M = [60.142611 + 6.100887 T - 2.709407 T^2 + 1.366779 T^3] R_E$$

where  $T = (t - t_{\text{now}}) / \text{Gyr}$ . This is a polynomial fit to over 50+ years of data. The

validity of the formula will end in 1 to 2 Gyr, when Earth's oceans

# Tides



Damped oscillator as a model equation of tidal wave

$$\ddot{x} + \underbrace{\gamma \dot{x}}_{\text{braking}} + \underbrace{\omega_0^2 x}_{\text{Free oscil.}} = \underbrace{\frac{F}{m} \cos \omega_p t}_{\text{forcing}}$$

$x =$  Deformation (vertical, relative to sphere)

$\omega_0 =$  Natural frequency of oscillations = (spring constant/m)<sup>1/2</sup>

$\omega_p =$  Frequency of forcing = spin frequency = angular speed of planet

$x$  = deformation w.r.t. sphere

$\omega_0$  = natural frequency of oscillation

$\omega_p$  = spin frequency of the planet

Damped harmonic oscillator solution:

$$x(t) = \underbrace{A e^{-\gamma t/2} \cos(\omega t + B)}_{\text{free oscillations}} + \underbrace{\frac{F/m}{\sqrt{\omega^4 + \gamma^2 \omega_p^2}} \cos(\omega_p t - \delta)}_{\text{tide}}$$

where  $\tan \delta = \gamma \omega_p / \omega^2$ ;  $\omega^2 = \omega_0^2 - \omega_p^2$

and  $A, B = \text{const.}$

tide amplitude  $\approx \frac{F}{m \omega_0^2}$ , since  $\omega_p^2 \ll \omega_0^2$

tidal lag angle  $\delta \approx \frac{\gamma \omega_p}{\omega_0^2}$

We can now compute  $E_0 = \underbrace{\left(\frac{F}{m \omega_0^2}\right) \cdot m \frac{\omega_0^2}{2}}_{\text{oscillator energy}}$

$$Q = \frac{E_0}{\left(\frac{2\pi}{\omega_p}\right) \oint \dot{E} dt} \cdot \omega_p$$

how many timescales =  $\omega_p^{-1}$ , before tides dampen? ?

$$|\dot{E}| = m \cdot \gamma \cdot \dot{x}^2 \quad (\text{force} \times \text{velocity})$$

Oscillator goodness parameter

$$Q \approx \frac{1}{\delta}$$

$\frac{F}{m} = \Delta f =$  force stretching the planet  
(per unit mass)

$$\Delta f = \frac{GM_S}{(r-r_p)^2} - \frac{GM_S}{r^2} \approx \frac{2GM_S r_p}{r^3}$$



$$\Delta M \sim \theta(L) \cdot M_S \cdot \left(\frac{r_p}{r}\right)^3$$

$$L = \frac{\delta}{2} r_p GM_S (\Delta M) \left[ \frac{1}{(r-p)^3} - \frac{1}{(r+p)^3} \right]$$

$$L \sim \frac{GM_S^2 r_p^5}{Q r^6} = \text{tidal torque}$$

acts on the planet, and on the satellite  
(signs opposite)

Consequences :

$\omega_p \searrow$  drops

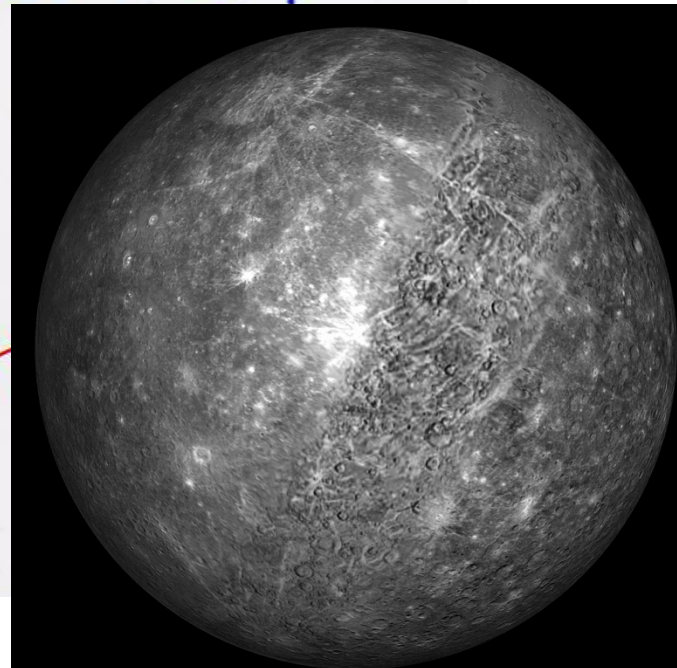
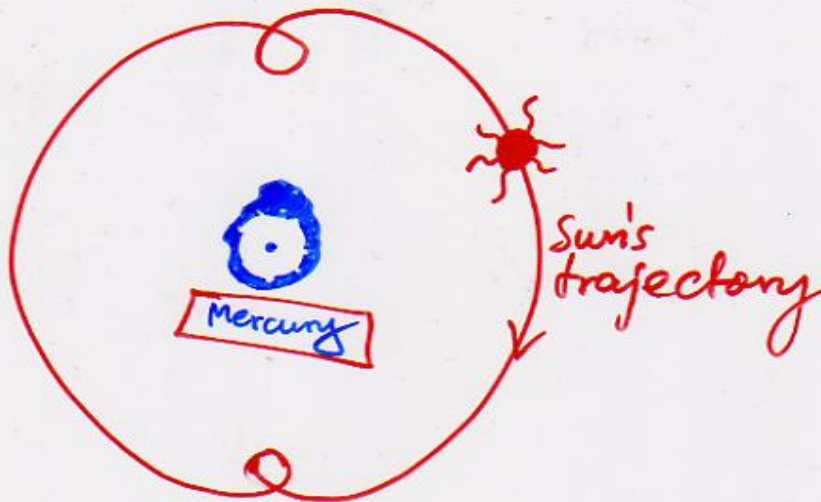
$r \nearrow$  grows

$\oplus + \ominus$
$0.0016 \frac{s}{centi.}$
$4 \frac{cm}{yr}$

Final state :  $\omega_p = 0$  ,  $\omega_p : n = 1:1$  ,  
planet in 1:1 corotation with  $M_S$

e.g. Mercury was thought  
& Moon was observed to corotate

However, Mercury is in  $\omega_p : n = 3:2$  !



Net torque = 0, stable resonance



(By coincidence, probably!) **Venus** rotates backwards  $\approx 5$  times between the approaches to Earth

( $Q = 1400 \Rightarrow$  that would be a physical reason...) but  $Q \ll 1400$ . E.g.  $\left\{ \begin{array}{l} Q \sim 15 \text{ for Earth} \\ Q \sim 100 \text{ without oceans} \end{array} \right.$

Q: Why does Venus rotate backwards?

A: 1) Giant impact; or

2) Massive atmosphere, atmospheric oscillations couple with insolation, which bloats the atmosphere

(There is such an effect, cancelling  $\sim 10\%$  of the gravitational tide, on Earth)



**Mars** :

Phobos  
Deimos

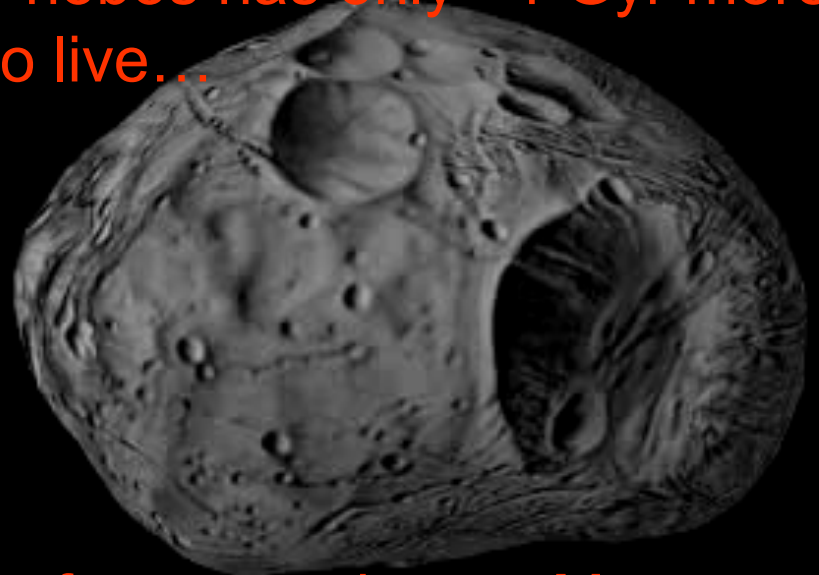
tidally locked spin  
like the Moon-Earth

Phobos very close, circles Mars  
3x faster than Mars spins  $\Rightarrow$   
 $\Rightarrow$   $r$   $\rightarrow$

this is a measurable effect :  $Q_{Mars} \sim 100$ ,



Phobos has only ~1 Gyr more to live...



before a crash onto Mars

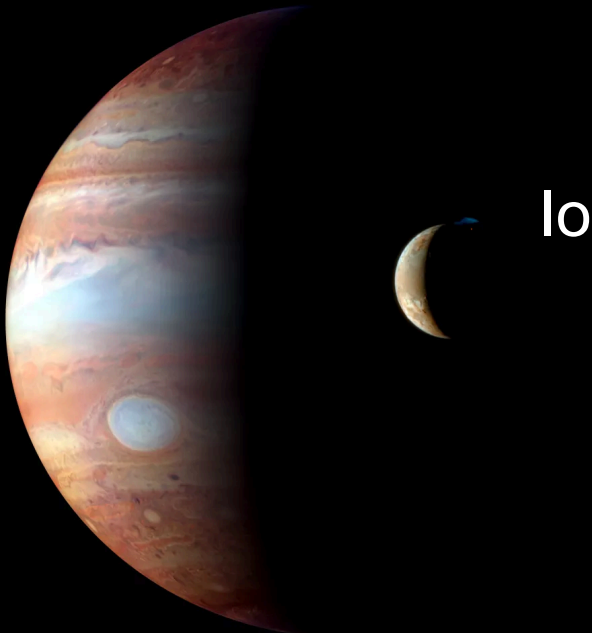
# Jupiter

All 4 Galilean satellites  
tidally locked in spin, as well.

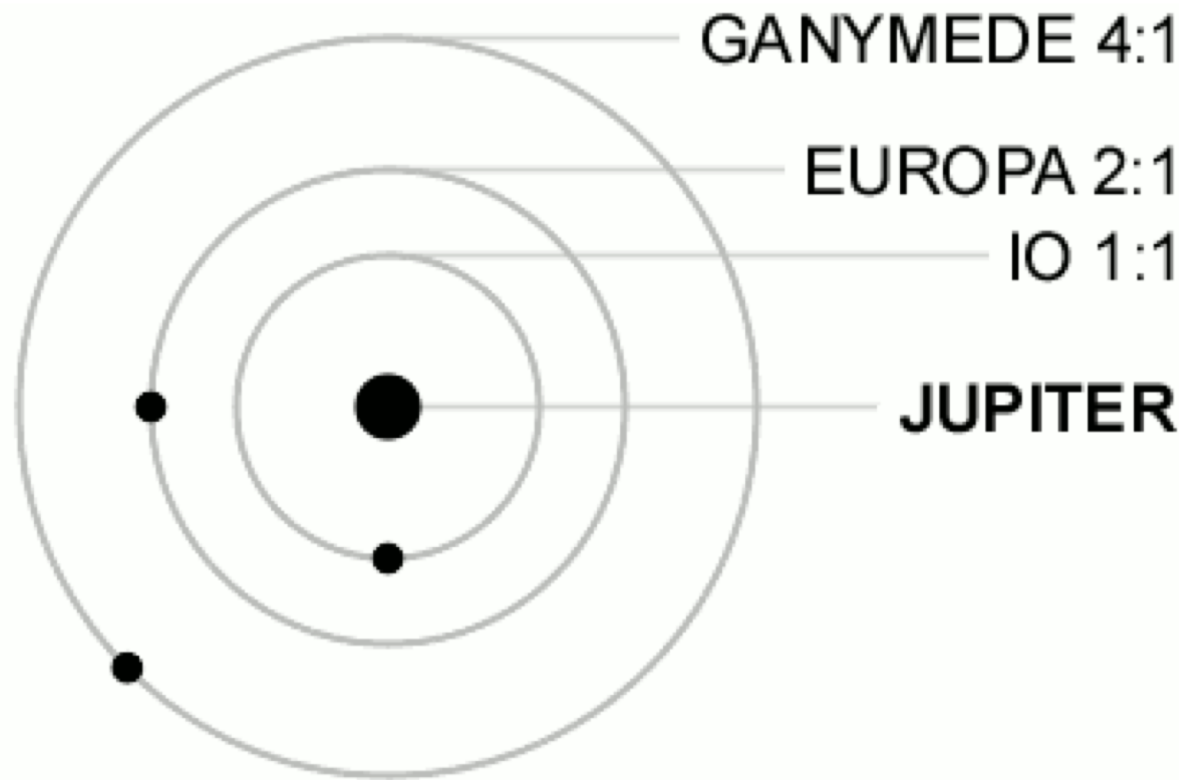
Q inside Jupiter must be very large  
&  $\delta$  very small or else these satellites  
would already have been pushed away!

$$Q \sim 6 \cdot 10^4$$

snapshot from New Horizons  
probe on the way to Pluto in 2007



Laplace resonance of Galilean satellites  
(except Callisto!) – the period ratios:

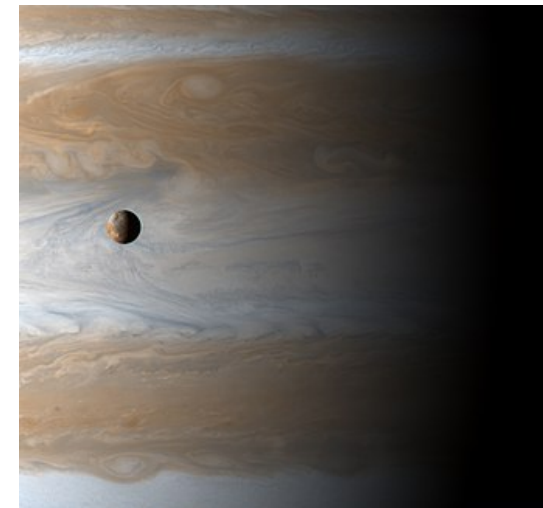


Physical radii

2634 km

1560 km

1821 km



Io

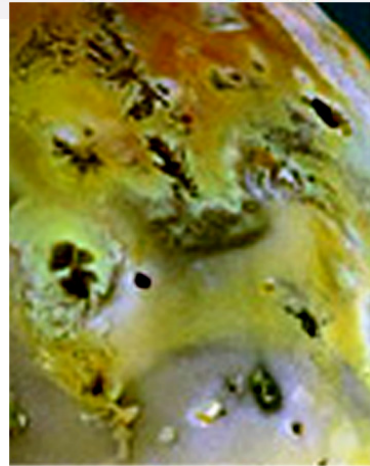
- tidally heated  $\Rightarrow$  the most active volcanoes in the Solar System!

Io - Europa 1:2 orbital resonance

forces eccentricity  $e_{Io} \approx 0.0041$

$\Rightarrow$  Io has nutation (oscillation) w.r.t. mean motion, Jupiter on Io's sky moves a little.

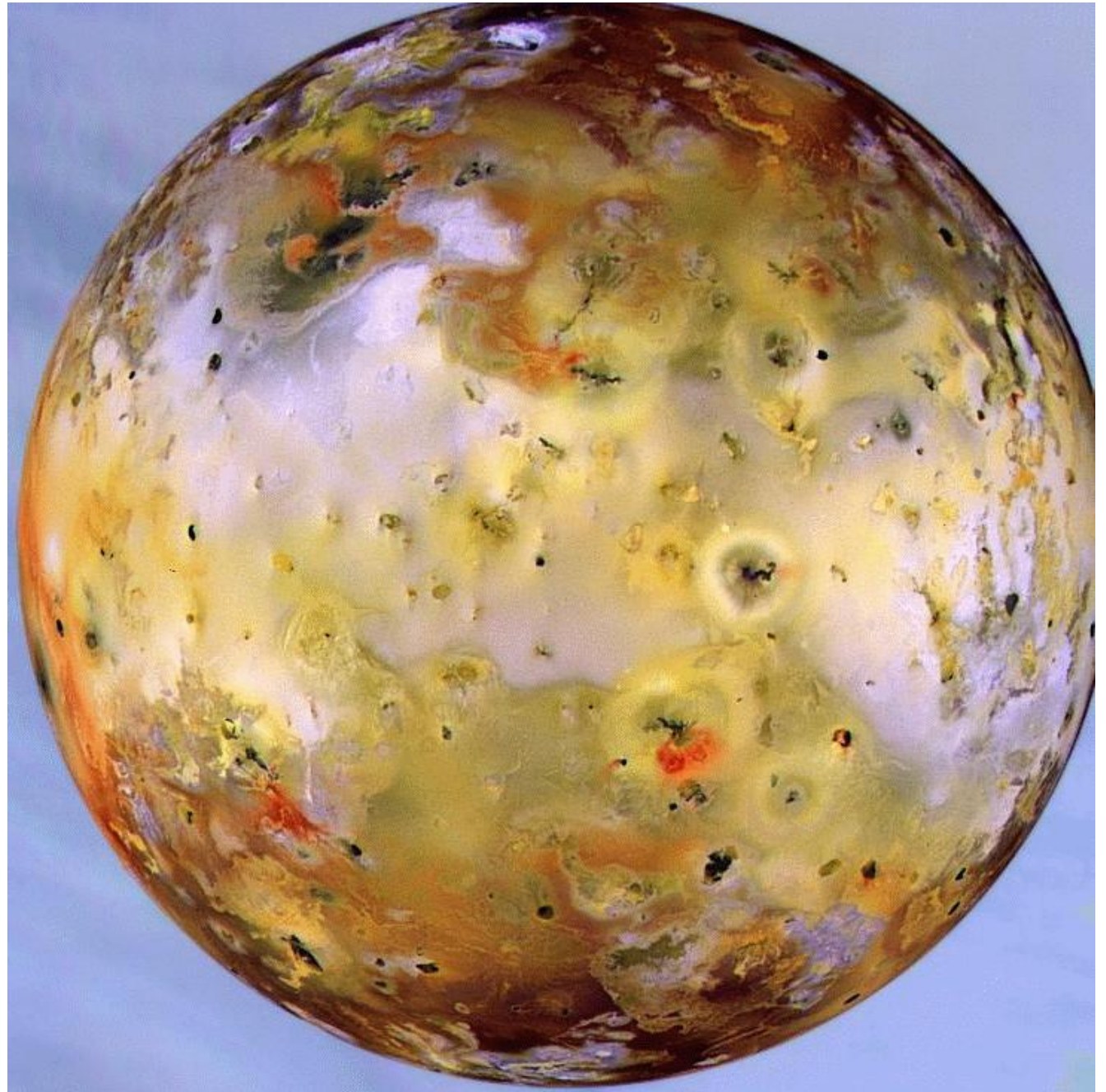
Io flexes, releases heat of friction, which melts the interior (predicted, then observed by Voyagers).



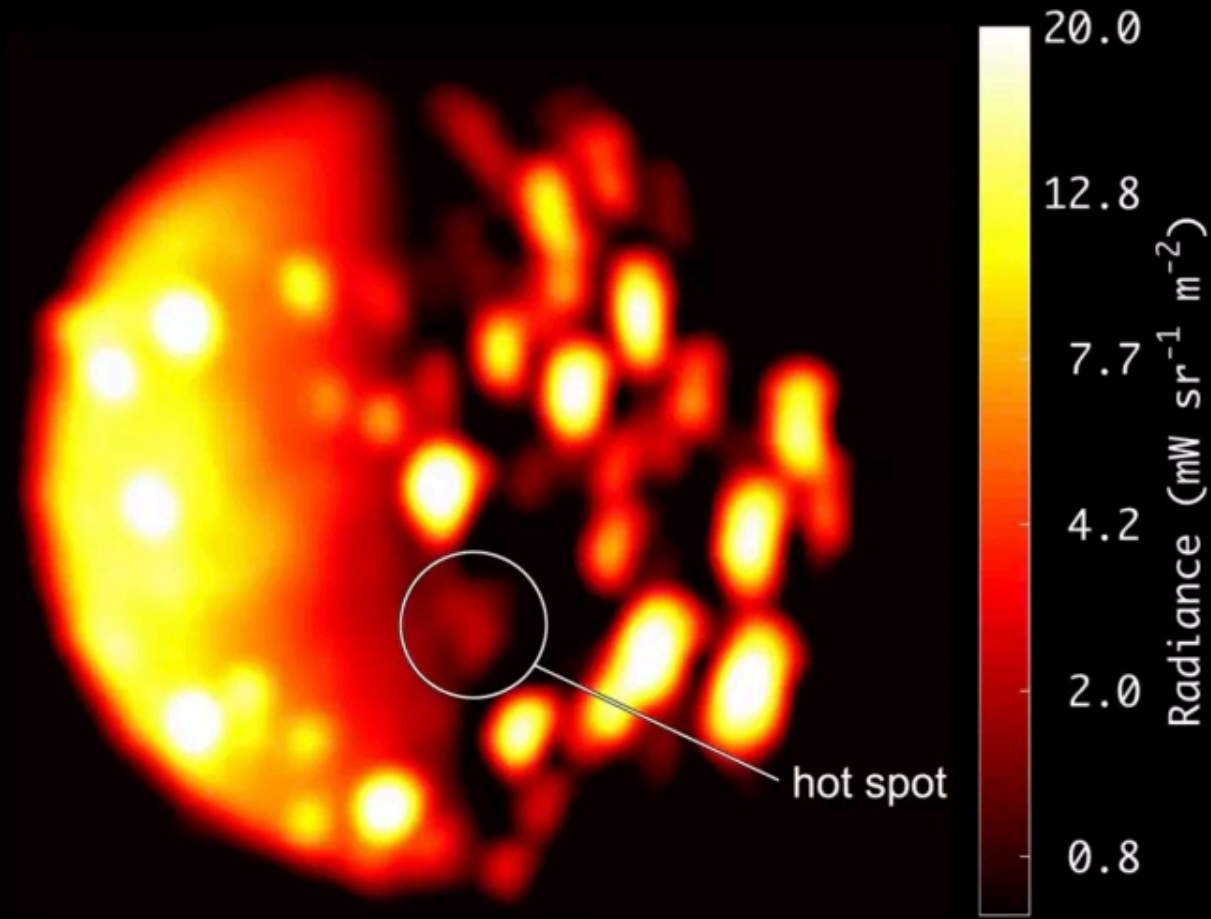
**Io** is tidally locked to Jupiter (same side faces it). Because of  $e > 0$  there is nutation.

Tidal flexing and heating of **Io** causes the most active volcanism in our system. It has 400 active volcanoes & produces  $10^{14}$  W of heat.

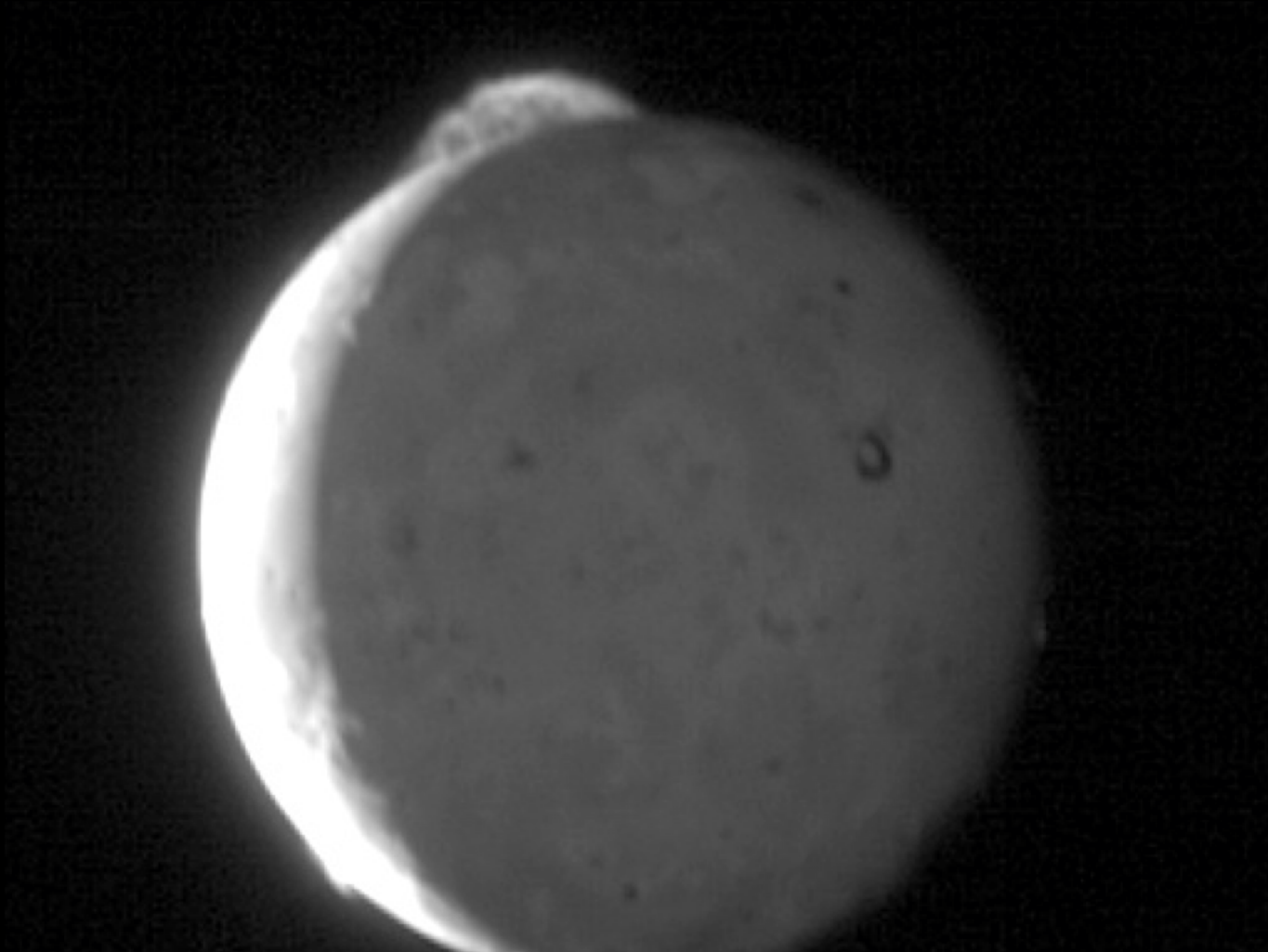
Also the moons Europa & Enceladus are tidally heated which allows underground water oceans on them.



Juno probe that left Earth in 2011 and since 2016 orbits Jupiter in an elliptic orbit took this snapshot of IR heat map of Io. Multiple volcanoes are visible. Juno's main mission ended in 2021.

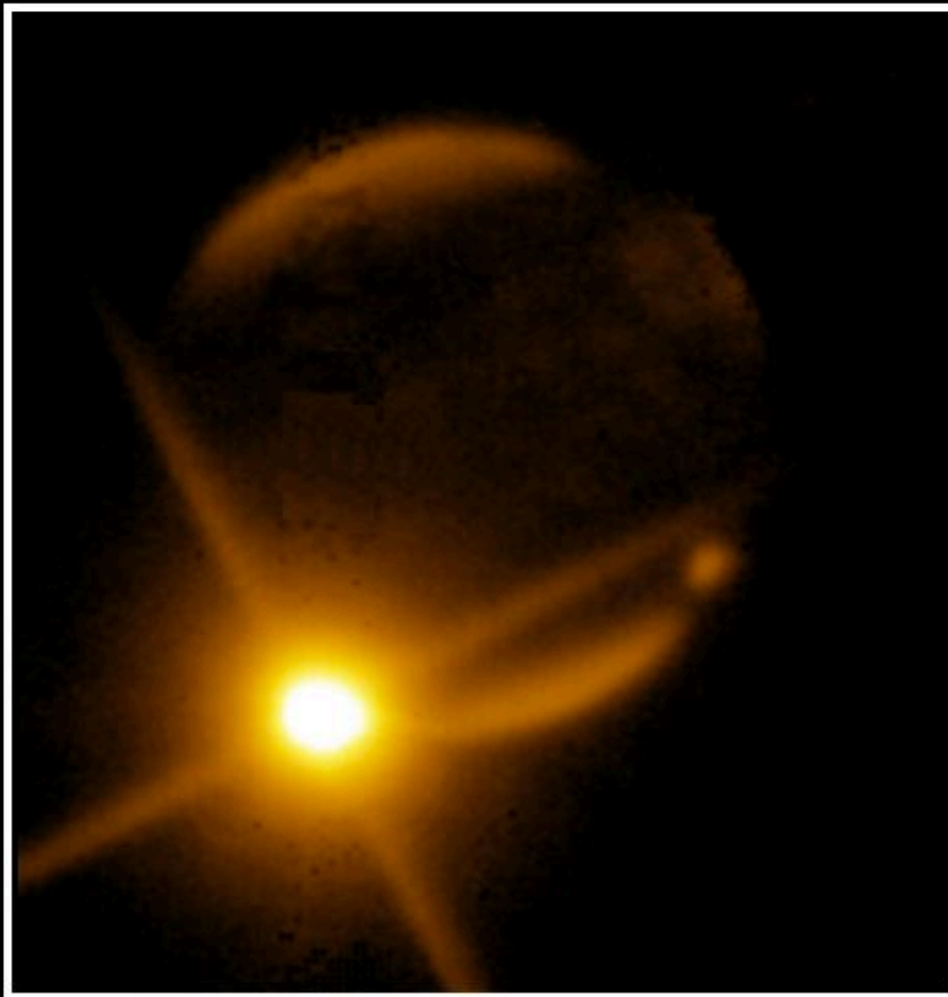


Plume from Tvashtar volcano on Io reaches 300 km height



spacecraft: New Horizons (NASA/JPL)



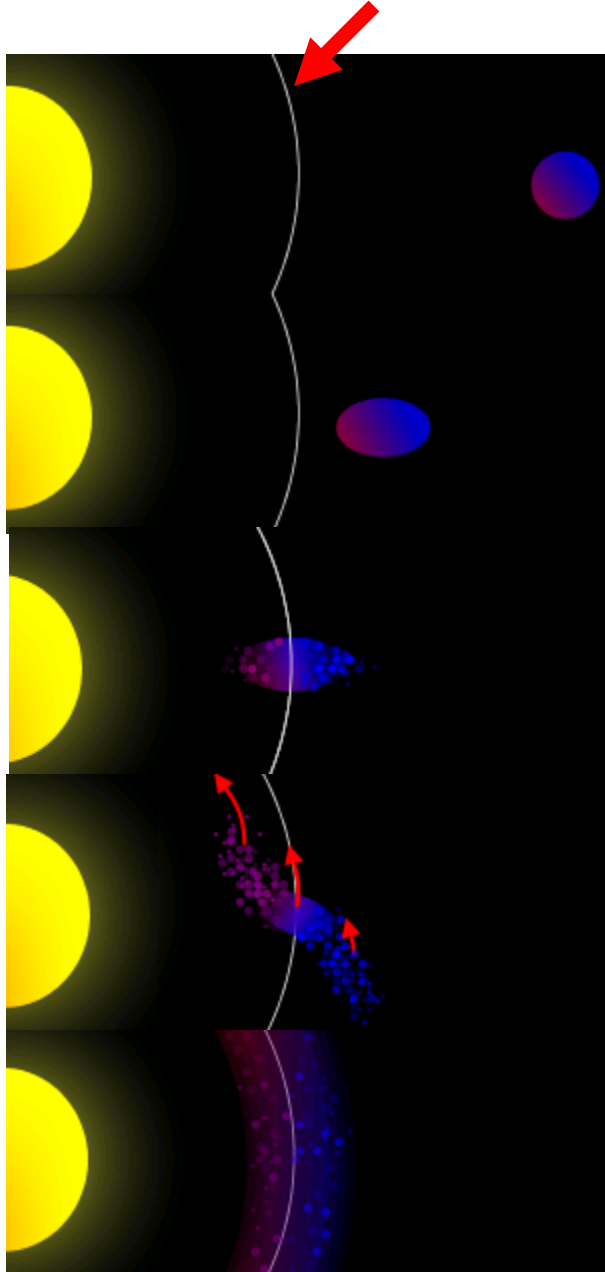


1994

**Impact of Fragment G of Comet Shoemaker-Levy on Jupiter  
The fireball is seen 12 minutes after impact at 2.34 microns.  
The impact A site is seen on the opposite limb of the planet.**

**Image at 2.34 microns with CASPIR by Peter McGregor  
ANU 2.3m telescope at Siding Spring**

**Roche limit:** tidal forces destroy bodies which come too close to a source of gravitation, through differential acceleration (tide)



Consider an orbiting mass of fluid held together by gravity. Far from the Roche limit the mass is practically spherical.

The body is deformed by tidal forces

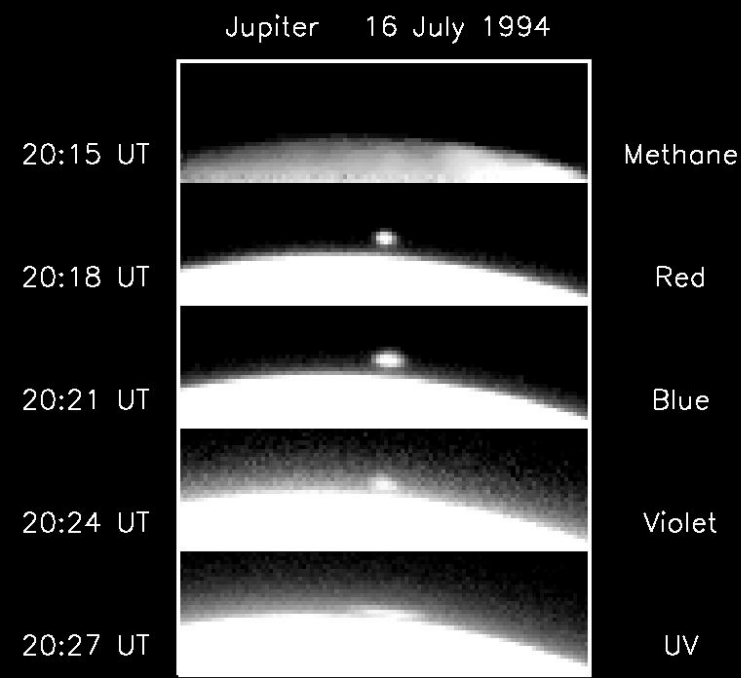
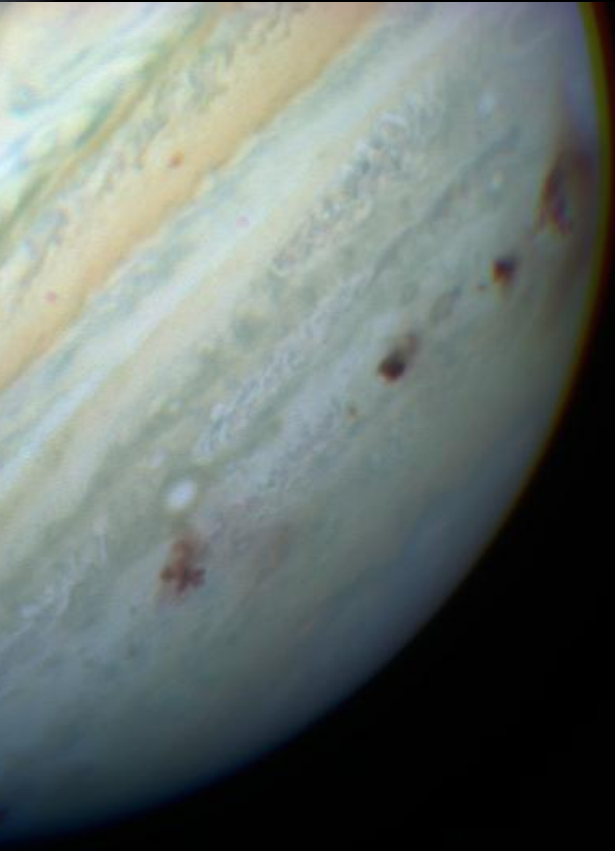
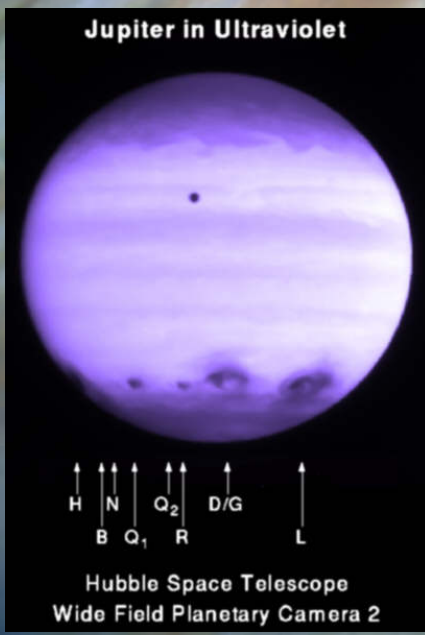
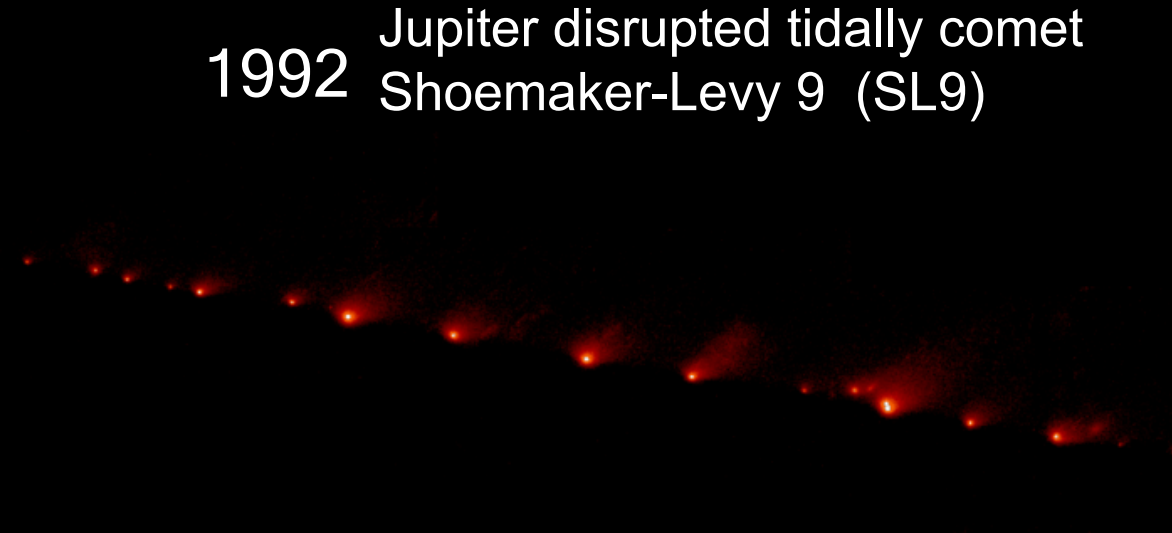
Within the Roche limit the mass' own gravity (self-gravity) can no longer withstand the tidal forces, and the body disintegrates.

Particles closer to the primary orbit move more quickly than particles farther away, as represented by the red arrows.

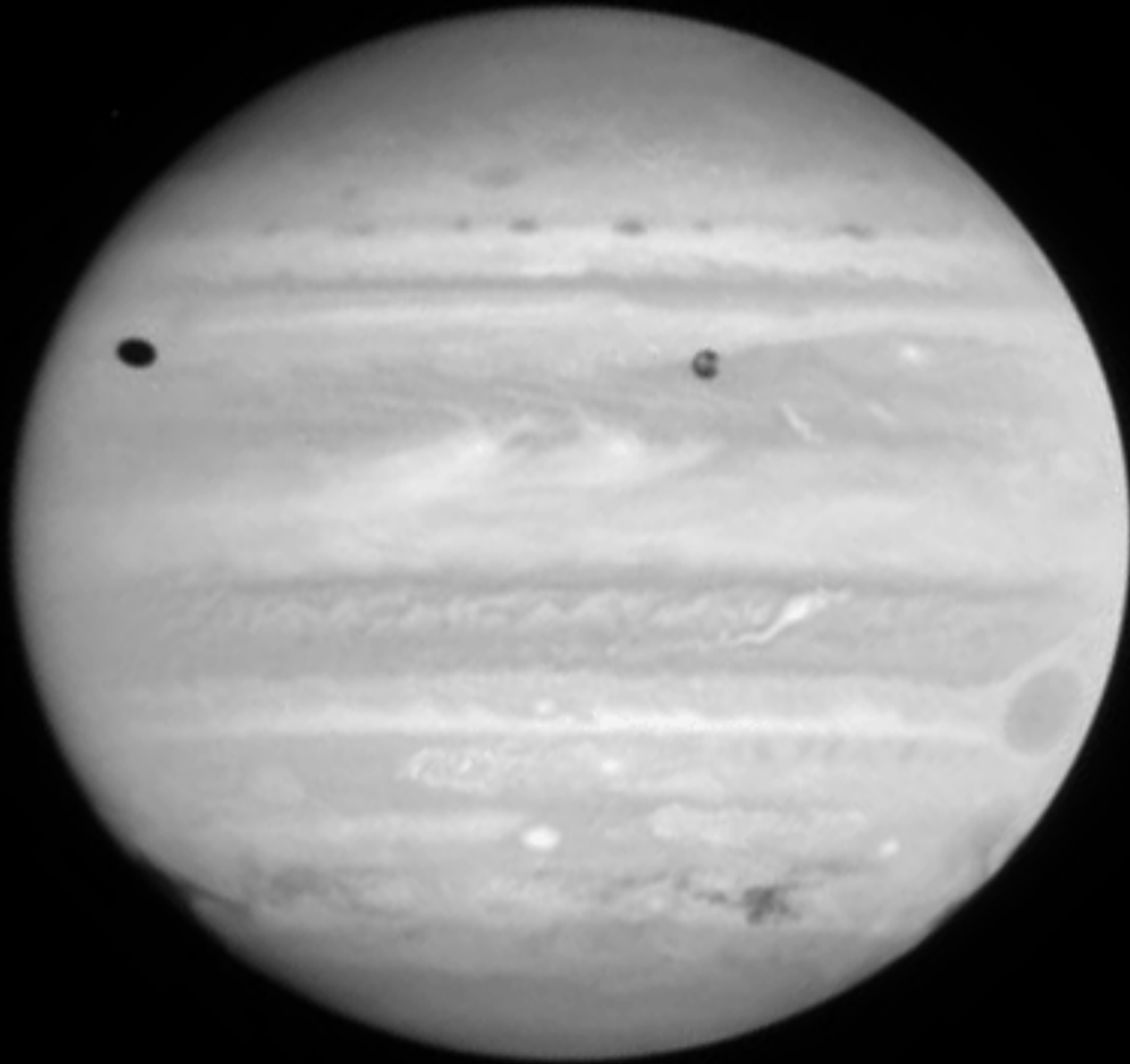
The varying orbital speed of the material eventually causes it to form a ring.

[http://www.answers.com/topic/roche-limit#wp-Fluid\\_satellites](http://www.answers.com/topic/roche-limit#wp-Fluid_satellites)

# 1992 Jupiter disrupted tidally comet Shoemaker-Levy 9 (SL9)



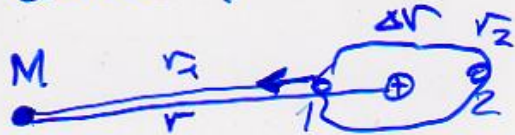
Wide Field Planetary Camera 2  
Hubble Space Telescope



Also see <https://www.youtube.com/watch?v=7zNuT4dbdjU>

First step in the computation is the estimation of the differential torque

tidal "force" = differential acceler.



Édouard Roche  
(1820–1883)

$$f_{\frac{1}{2}} = -\frac{GM}{r_{\frac{1}{2}}^2}$$

$$\Delta f = f_2 - f_1 = -GM \left( \frac{1}{r_2^2} - \frac{1}{r_1^2} \right)$$

$$\Delta r = r_2 - r_1$$

$$\Delta f = -\frac{GM}{r_1^2} \left( \frac{r_1^2}{(r_1 + \Delta r)^2} - 1 \right)$$

$$\left[ 1 + \left( \frac{\Delta r}{r_1} \right) \right]^{-2} \approx 1 - 2 \frac{\Delta r}{r_1} \ll 1$$

$$\Delta f \approx + \frac{2GM}{r_1^2} \cdot \frac{\Delta r}{r_1} \ll 1$$

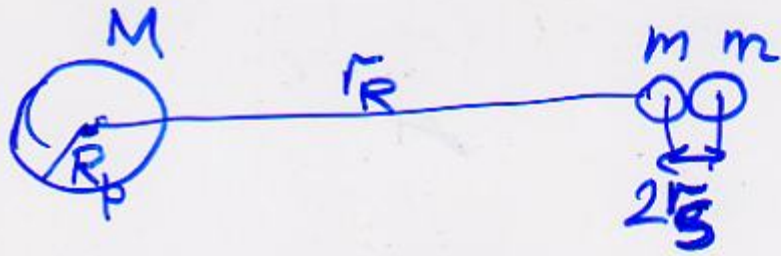
**Roche limit:**

*Fluid limit:*  
tides destroy a body bound by self-gravity within that minimum distance  $r_R$

*Solid/rocky limit:*  
the same but with additional binding (electromagnetic) forces of material strength

Q: which limit is larger?

# Roche limit calculation



$$\Delta f = \frac{2GM}{r_R^2} \cdot \frac{2r_s}{r_R}$$

tidal "force"

$$f_{gr} = \frac{Gm}{(2r_s)^2}$$

gravity

$$\frac{4M r_s}{r_R^3} = \frac{m}{4r_s^2}$$

$$16 \frac{M}{m} = \frac{r_R^3}{r_s^3}$$

$$\frac{M}{m} = \left( \frac{R_p^3 \rho_p}{r_s^3 \rho_s} \right)$$

$$16 \frac{\rho_p}{\rho_s} = \frac{r_R^3}{r_p^3}$$

$$\Rightarrow r_R = r_p \left( \frac{\rho_p}{\rho_s} \right)^{\frac{1}{3}}$$

2.55

$$\rightarrow 4 \frac{M}{R^3} r_s = \frac{M}{4 r_s^2}$$

$$16 \frac{M}{m} = \frac{R^3}{r_s^3}$$

$$\frac{M}{m} = \left( \frac{R_p^3 \rho_p}{r_s^3 \rho_s} \right)$$

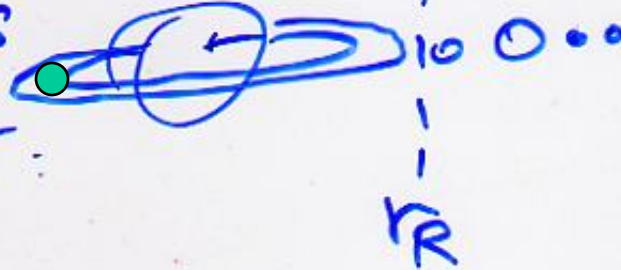
$$16 \frac{\rho_p}{\rho_s} = \frac{R^3}{r_s^3} \Rightarrow r_R = R_p \left( \frac{\rho_p}{\rho_s} \right)^{\frac{1}{3}}$$

2.55

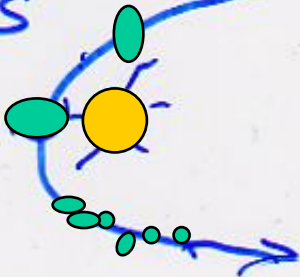
Roche limit

[pron.: Rosh]

Applications:  
 ① satellite systems



② comets



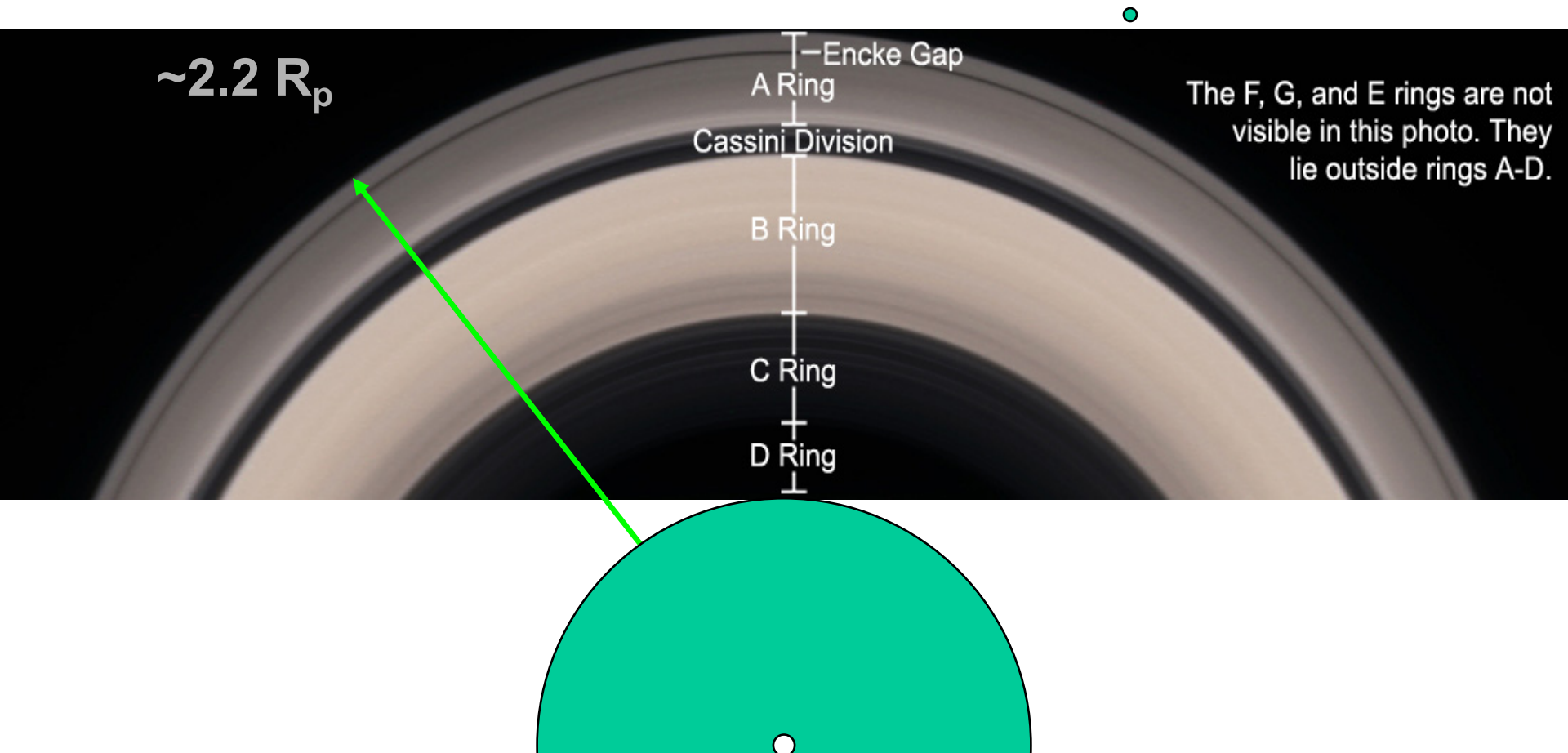
Comet Shoemaker-Levi SL-9 (1994) disintegrated after approach to Jupiter. So do the sun-grazing comets

Roche limit for fluid satellite (1849):  $d \approx 2.44R \left( \frac{\rho_M}{\rho_m} \right)^{1/3}$

Roche limit for rocky satellite:  $d \approx 1.4R \left( \frac{\rho_M}{\rho_m} \right)^{1/3}$

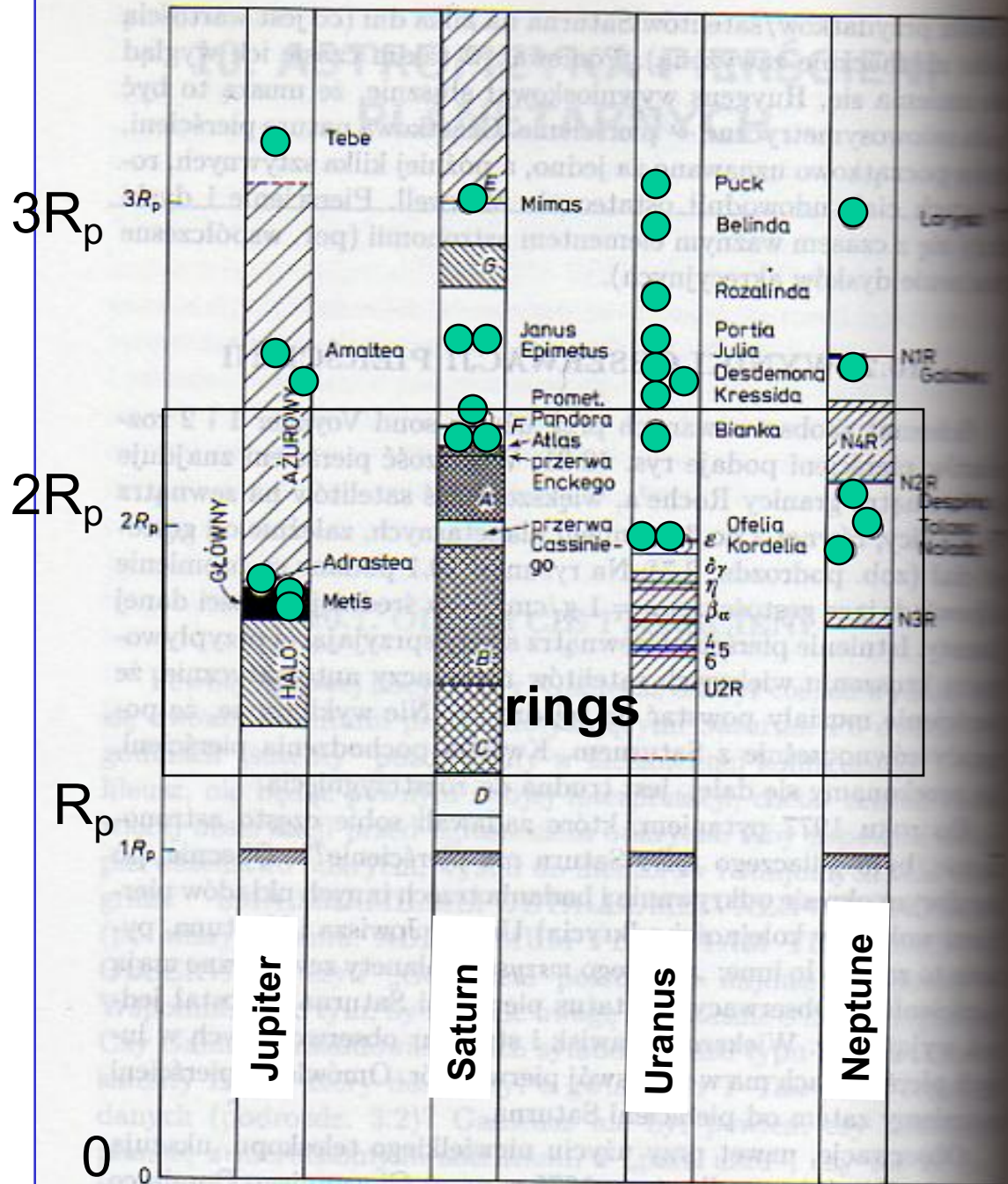
( $R_p$  = planet's radius,  $\rho_M$  = its density)

→ **satellites generally are found outside the rings**





● = satellite



**Roche limit** is somewhere between  $1.4$  and  $2 R_p$ : satellites are destroyed if found closer to their planets.

*Ring systems drawn to scale of their planets, but not to relative scale.*

# Lecture L06 - ASTC25

## Elements of Celestial Mechanics

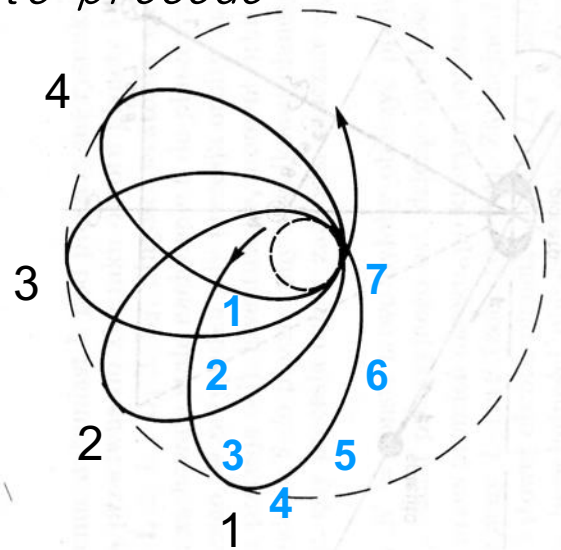
Precession of orbits and rotational axes

Lunisolar precession of Earth's spin axis

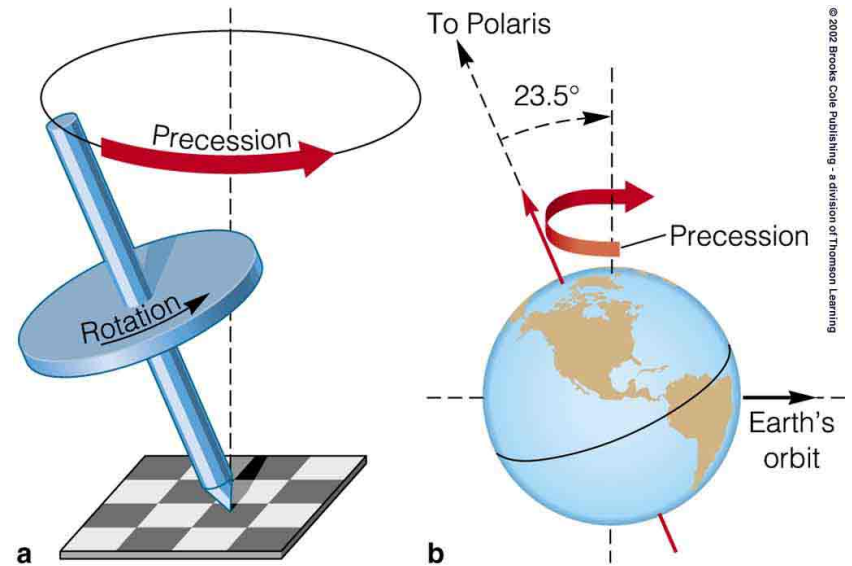
# Two kinds of Precession

orbit precession,

*Praecedere [latin] =  
to precede*



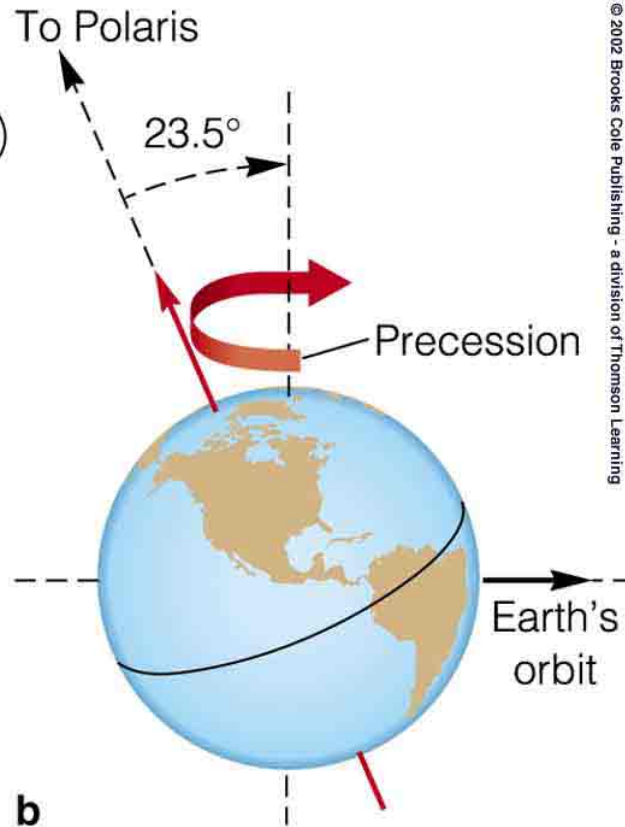
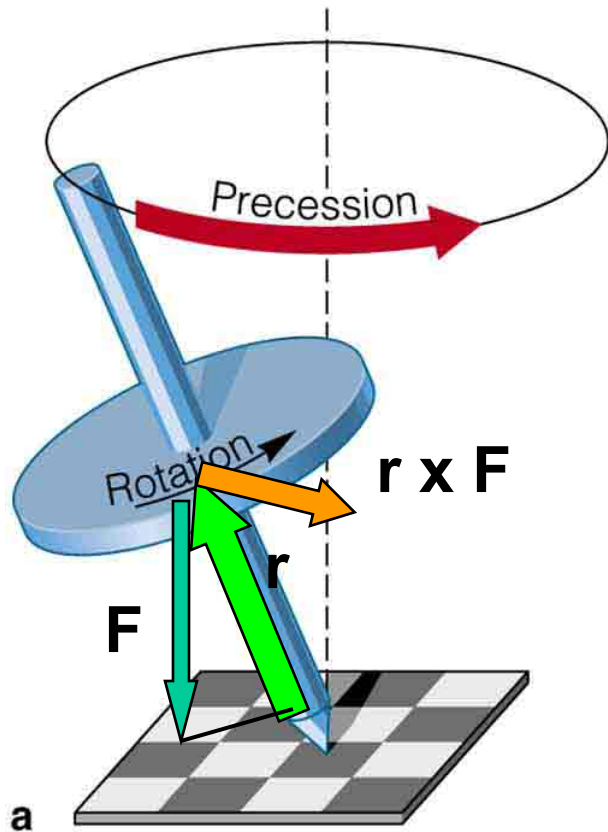
spin axis precession



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Exceptions (closed orbits):  
two force laws only:  $F \sim r^{-2}$ , and  $F \sim r$   
(Bertrand's theorem)

Angular momentum vector shifted sideways by torque, value preserved  $\Rightarrow$  precession.  $\frac{d\vec{L}}{dt} = (\text{torque})$   
 $\frac{d\vec{L}}{dt} = (\vec{r} \times \vec{F}) dt$

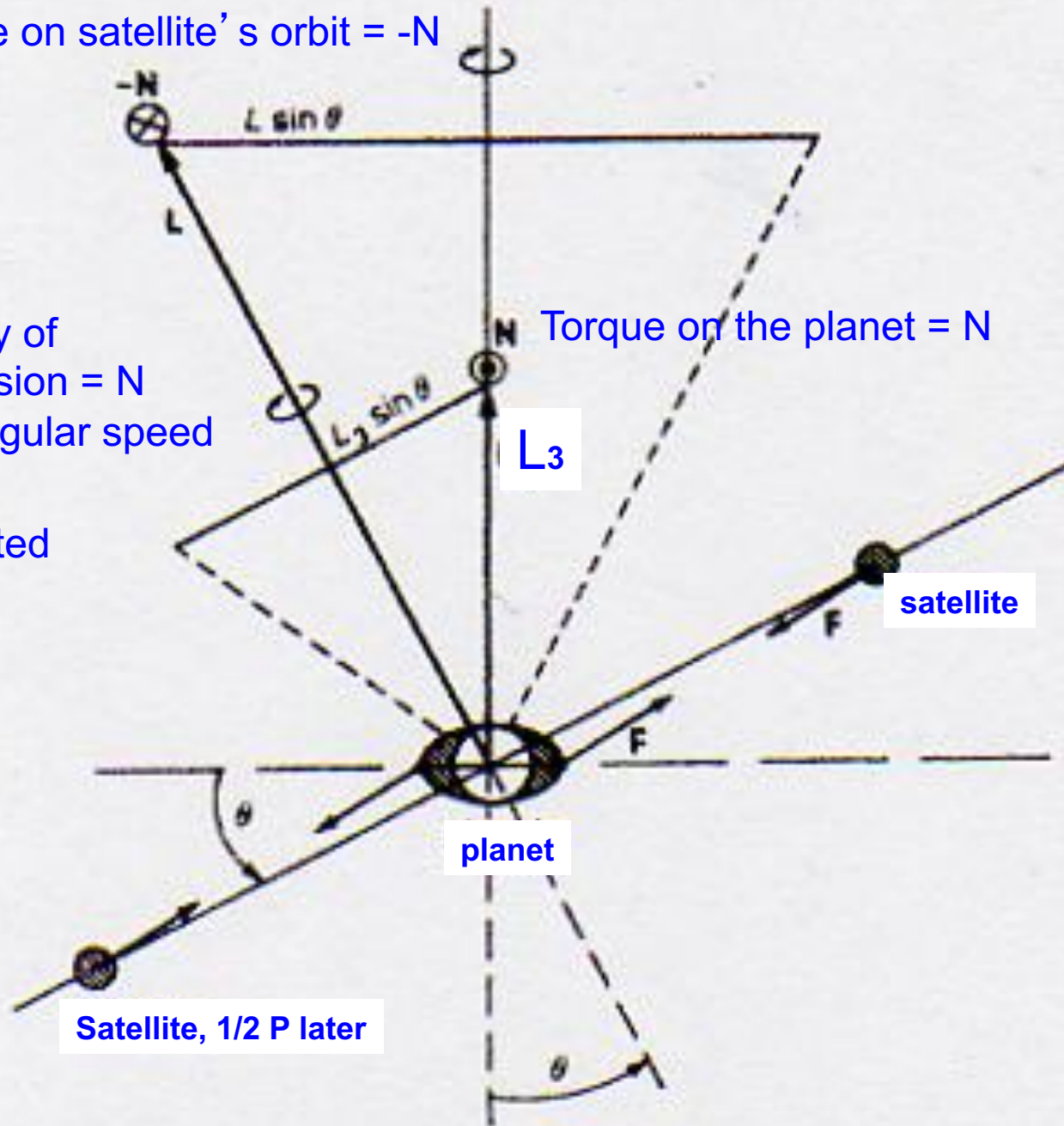


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Torque on satellite's orbit =  $-N$

Velocity of precession =  $N$   
 $\implies$  angular speed can be computed



$r$  = perturber distance

perturber mass

Obliquity (inclination)

$$N = -\frac{3GM_K(I_3 - I_1)}{2r^3} \sin \theta \cos \theta,$$

**Earth axis' precession (angular speed):**

$$\omega_p = \frac{N}{\omega_3 I_3 \sin \theta}$$

$$\omega_p = -\frac{3GM_K}{2\omega_3 r^3} \frac{I_3 - I_1}{I_3} \cos \theta.$$

*$I_{1,2,3}$  = three moments of inertia of the Earth, satisfying:*

$$\frac{I_3 - I_1}{I_3} = 0,00327,$$

*(from Earth flattening)*

**The corresponding period of precession of Earth axis = 33000 yr (from the Moon), and 81000 (from the sun)**

**Combined *luni-solar precession* has period 26000 yr = 51 "/yr**  
(Note:  $1/33000 + 1/81000 = 1/26000$ , angular speeds, not periods, add up).  
**Precession affects right ascension of objects in catalogues and maps, which therefore must state the "epoch" of coordinates.**